

# The antifield-BRST approach to (gauge) field theories: an overview

Marc Henneaux

Exceptional Dimensions  
A Conference in Honour of Bernard Julia  
(16-17 December 2019, Institut Henri Poincaré)

# Introduction

## The antifield-BRST approach to (gauge) field theories: an overview

Marc Henneaux

### Introduction

BRST differential in Yang-Mills theory

Antifields and Koszul-Tate resolution

Homological perturbation theory

Local functionals and characteristic complex

Conclusions

# Introduction

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

## Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

Bernard has always emphasized the crucial importance of cohomological ideas in physics.

# Introduction

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

**Introduction**

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

Bernard has always emphasized the crucial importance of cohomological ideas in physics.

**My talk will be an illustration of the power of cohomological methods in the context of gauge field theories.**

# Introduction

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

**Introduction**

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

Bernard has always emphasized the crucial importance of cohomological ideas in physics.

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It will be devoted to BRST theory.

# Introduction

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

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I will focus in particular on the cohomological significance of the antifields,

# Introduction

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

Bernard has always emphasized the crucial importance of cohomological ideas in physics.

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which is crucial for computing explicitly the BRST cohomology.

# Introduction

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

Bernard has always emphasized the crucial importance of cohomological ideas in physics.

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which is crucial for computing explicitly the BRST cohomology.

These were introduced by Zinn-Justin, and Batalin and Vilkovisky.



# Introduction

## The antifield-BRST approach to (gauge) field theories: an overview

Marc Henneaux

### Introduction

BRST differential in Yang-Mills theory

Antifields and Koszul-Tate resolution

Homological perturbation theory

Local functionals and characteristic complex

Conclusions

I will successively discuss :

# Introduction

## The antifield-BRST approach to (gauge) field theories: an overview

Marc Henneaux

### Introduction

BRST differential in Yang-Mills theory

Antifields and Koszul-Tate resolution

Homological perturbation theory

Local functionals and characteristic complex

Conclusions

I will successively discuss :

- **BRST differential in Yang-Mills theory**

# Introduction

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

I will successively discuss :

- BRST differential in Yang-Mills theory
- Antifields and Koszul-Tate resolution

# Introduction

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

## Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

I will successively discuss :

- BRST differential in Yang-Mills theory
- Antifields and Koszul-Tate resolution
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# Introduction

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

I will successively discuss :

- BRST differential in Yang-Mills theory
- Antifields and Koszul-Tate resolution
- Homological perturbation theory (and  $L_\infty$  algebras)
- Local functionals and characteristic complex

# Introduction

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

## Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

I will successively discuss :

- BRST differential in Yang-Mills theory
- Antifields and Koszul-Tate resolution
- Homological perturbation theory (and  $L_\infty$  algebras)
- Local functionals and characteristic complex
- Conclusions

# Yang-Mills theory

## The antifield-BRST approach to (gauge) field theories: an overview

Marc Henneaux

Introduction

**BRST differential in Yang-Mills theory**

Antifields and Koszul-Tate resolution

Homological perturbation theory

Local functionals and characteristic complex

Conclusions

# Yang-Mills theory

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

**BRST differential  
in Yang-Mills  
theory**

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

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“minimal sector”)



# Yang-Mills theory

The antifield-BRST approach to (gauge) field theories: an overview

Marc Henneaux

Introduction

BRST differential in Yang-Mills theory

Antifields and Koszul-Tate resolution

Homological perturbation theory

Local functionals and characteristic complex

Conclusions

The BRST differential  $s$  in Yang-Mills theory reads (in the “minimal sector”)

$$sA_\mu^a = D_\mu C^a, \quad sC^a = -\frac{1}{2}f_{bc}^a C^b C^c, \\ sA_a^{*\mu} = D_\nu F_a^{\nu\mu} + f_{ac}^b A_b^{*\mu} C^c, \quad sC_a^* = D_\mu A_a^{*\mu} + f_{ac}^b C_b^* C^c.$$

# Yang-Mills theory

The antifield-BRST approach to (gauge) field theories: an overview

Marc Henneaux

Introduction

BRST differential in Yang-Mills theory

Antifields and Koszul-Tate resolution

Homological perturbation theory

Local functionals and characteristic complex

Conclusions

The BRST differential  $s$  in Yang-Mills theory reads (in the “minimal sector”)

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It is generated in the antibracket by the solution  $S$  of the “master equation”,

# Yang-Mills theory

The antifield-BRST approach to (gauge) field theories: an overview

Marc Henneaux

Introduction

BRST differential in Yang-Mills theory

Antifields and Koszul-Tate resolution

Homological perturbation theory

Local functionals and characteristic complex

Conclusions

The BRST differential  $s$  in Yang-Mills theory reads (in the “minimal sector”)

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It is generated in the antibracket by the solution  $S$  of the “master equation”,

$$sF = (S, F) \\ S = -\frac{1}{4} \int d^n x F_a^{\mu\nu} F_{\mu\nu}^a + \int d^n x A_a^{*\mu} sA_\mu^a + \int d^n x C_a^* sC^a,$$

# Yang-Mills theory

The antifield-BRST approach to (gauge) field theories: an overview

Marc Henneaux

Introduction

BRST differential in Yang-Mills theory

Antifields and Koszul-Tate resolution

Homological perturbation theory

Local functionals and characteristic complex

Conclusions

The BRST differential  $s$  in Yang-Mills theory reads (in the “minimal sector”)

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# Yang-Mills theory

The antifield-BRST approach to (gauge) field theories: an overview

Marc Henneaux

Introduction

BRST differential in Yang-Mills theory

Antifields and Koszul-Tate resolution

Homological perturbation theory

Local functionals and characteristic complex

Conclusions

The BRST differential  $s$  in Yang-Mills theory reads (in the “minimal sector”)

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$$s^2 = 0 \Leftrightarrow (S, S) = 0.$$

# Antifields

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

**BRST differential  
in Yang-Mills  
theory**

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

# Antifields

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

**BRST differential  
in Yang-Mills  
theory**

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

$A_a^{*\mu}$  and  $C_a^*$  are the “antifields”.

# Antifields

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

$A_a^{*\mu}$  and  $C_a^*$  are the “antifields”.

Antifields were originally introduced by Zinn-Justin in his seminal work on the renormalization of gauge theories, as sources coupled to the BRST variations of the fields.



# Antifields

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

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This was motivated by the desire to control how the nonlinear BRST symmetry passes through the renormalization process.

# Antifields

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

$A_a^{*\mu}$  and  $C_a^*$  are the “antifields”.

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A different interpretation of the antifields can be developed.

# Antifields

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

$A_a^{*\mu}$  and  $C_a^*$  are the “antifields”.

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This interpretation has cohomological origins and views the antifields as the generators of a differential complex that implements the gauge invariant equations of motion in cohomology.

# Antifields

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

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A different interpretation of the antifields can be developed.

This interpretation has cohomological origins and views the antifields as the generators of a differential complex that implements the gauge invariant equations of motion in cohomology.

This different point of view turns out to be crucial for computing the BRST cohomology.

# Covariant phase space

## The antifield-BRST approach to (gauge) field theories: an overview

Marc Henneaux

Introduction

BRST differential in Yang-Mills theory

**Antifields and Koszul-Tate resolution**

Homological perturbation theory

Local functionals and characteristic complex

Conclusions

# Covariant phase space

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

The phase space  $\Pi$  of a gauge theory can be covariantly described as the space of solutions to the equations of motion modulo the gauge transformations.

# Covariant phase space

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

The phase space  $\Pi$  of a gauge theory can be covariantly described as the space of solutions to the equations of motion modulo the gauge transformations.

The equations of motion define a “surface” in the space  $J$  of all histories, which is called the “stationary surface” and denoted by  $\Sigma$ .

# Covariant phase space

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

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$C^\infty(\Sigma)$  is the space of smooth functions on that surface.



# Covariant phase space

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

The phase space  $\Pi$  of a gauge theory can be covariantly described as the space of solutions to the equations of motion modulo the gauge transformations.

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$C^\infty(\Sigma)$  is the space of smooth functions on that surface.

Formally,  $\Pi$  is the quotient space  $\Pi = \Sigma/\mathcal{O}$  of the stationary surface  $\Sigma$  by the gauge orbits  $\mathcal{O}$  generated by the gauge transformations.

# Covariant phase space

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

The phase space  $\Pi$  of a gauge theory can be covariantly described as the space of solutions to the equations of motion modulo the gauge transformations.

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[For local objects, jet space formalism can be used to put these considerations on a firmer footing.]

# Covariant phase space

## The antifield-BRST approach to (gauge) field theories: an overview

Marc Henneaux

Introduction

BRST differential in Yang-Mills theory

**Antifields and Koszul-Tate resolution**

Homological perturbation theory

Local functionals and characteristic complex

Conclusions

# Covariant phase space

## The antifield-BRST approach to (gauge) field theories: an overview

Marc Henneaux

Introduction

BRST differential in Yang-Mills theory

**Antifields and Koszul-Tate resolution**

Homological perturbation theory

Local functionals and characteristic complex

Conclusions

The observables are the functions on  $\Pi$ .

# Covariant phase space

## The antifield-BRST approach to (gauge) field theories: an overview

Marc Henneaux

Introduction

BRST differential in Yang-Mills theory

Antifields and Koszul-Tate resolution

Homological perturbation theory

Local functionals and characteristic complex

Conclusions

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This description of the observables involves two steps :

# Covariant phase space

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

The observables are the functions on  $\Pi$ .

This description of the observables involves two steps :

(1) Restriction to the stationary surface ;

# Covariant phase space

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

The observables are the functions on  $\Pi$ .

This description of the observables involves two steps :

- (1) Restriction to the stationary surface ;
- (2) Implementation of the gauge invariance condition on  $\Sigma$ .

# Covariant phase space

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

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The BRST differential provides a cohomological formulation of  $C^\infty(\Pi)$  at ghost number zero,  $H^0(s) = \{Observables\}$ .



# Covariant phase space

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

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To each of the steps (1), (2) corresponds a separate differential.

# Covariant phase space

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

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To each of the steps (1), (2) corresponds a separate differential.

Both differentials appear in  $s$ .

# Antifield number

## The antifield-BRST approach to (gauge) field theories: an overview

Marc Henneaux

Introduction

BRST differential in Yang-Mills theory

**Antifields and Koszul-Tate resolution**

Homological perturbation theory

Local functionals and characteristic complex

Conclusions

# Antifield number

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

**Antifields and  
Koszul-Tate  
resolution**

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

To exhibit this property, it is useful to introduce the antifield number,

# Antifield number

The antifield-BRST approach to (gauge) field theories: an overview

Marc Henneaux

Introduction

BRST differential in Yang-Mills theory

Antifields and Koszul-Tate resolution

Homological perturbation theory

Local functionals and characteristic complex

Conclusions

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	puregh	antifd	gh
$A_\mu^a$	0	0	0
$C^a$	1	0	1
$A_a^{*\mu}$	0	1	-1
$C_a^*$	0	2	-2

Pure ghost number, antifield number and  $gh \equiv \text{puregh} - \text{antifd}$  ("total ghost number"), for the different field types

# Antifield number

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

To exhibit this property, it is useful to introduce the antifield number,

	puregh	antifd	gh
$A_\mu^a$	0	0	0
$C^a$	1	0	1
$A_a^{*\mu}$	0	1	-1
$C_a^*$	0	2	-2

Pure ghost number, antifield number and  $gh \equiv \text{puregh} - \text{antifd}$  ("total ghost number"), for the different field types

One has  $s = \delta + \gamma$ , with  $\text{antifd}(\delta) = -1$  and  $\text{antifd}(\gamma) = 0$

# Antifield number

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

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Explicitly,  $\delta A_\mu^a = 0$ ,  $\delta C^a = 0$ ,  $\delta A_a^{*\mu} = D_\nu F_a^{\nu\mu}$ ,  $\delta C_a^* = D_\mu A_a^{*\mu}$

# Antifield number

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

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and  $\gamma A_\mu^a = D_\mu C^a$ ,  $\gamma C^a = -\frac{1}{2} f_{bc}^a C^b C^c$ ,  $\gamma A_a^{*\mu} = f_{ac}^b A_b^{*\mu} C^c$ ,  
 $\gamma C_a^* = f_{ac}^b C_b^* C^c$ .



# Antifield number

The antifield-BRST approach to (gauge) field theories: an overview

Marc Henneaux

Introduction

BRST differential in Yang-Mills theory

Antifields and Koszul-Tate resolution

Homological perturbation theory

Local functionals and characteristic complex

Conclusions

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Nilpotency of  $s$  implies  $\delta^2 = 0$ ,  $\delta\gamma + \gamma\delta = 0$ ,  $\gamma^2 = 0$ .

# Koszul-Tate differential

## The antifield-BRST approach to (gauge) field theories: an overview

Marc Henneaux

Introduction

BRST differential in Yang-Mills theory

**Antifields and Koszul-Tate resolution**

Homological perturbation theory

Local functionals and characteristic complex

Conclusions

# Koszul-Tate differential

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

**Antifields and  
Koszul-Tate  
resolution**

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

The differential  $\delta$  is called the “Koszul-Tate differential” because it is associated with the Koszul-Tate resolution of the algebra of functions on the stationary surface (first step),

# Koszul-Tate differential

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

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# Koszul-Tate differential

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

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# Koszul-Tate differential

The antifield-BRST approach to (gauge) field theories: an overview

Marc Henneaux

Introduction

BRST differential in Yang-Mills theory

Antifields and Koszul-Tate resolution

Homological perturbation theory

Local functionals and characteristic complex

Conclusions

The differential  $\delta$  is called the “Koszul-Tate differential” because it is associated with the Koszul-Tate resolution of the algebra of functions on the stationary surface (first step),

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This second aspect is well appreciated (Chevalley-Eilenberg differential and “Lie algebra cohomology” in the relevant representation space).

# Koszul-Tate differential

The antifield-BRST approach to (gauge) field theories: an overview

Marc Henneaux

Introduction

BRST differential in Yang-Mills theory

Antifields and Koszul-Tate resolution

Homological perturbation theory

Local functionals and characteristic complex

Conclusions

The differential  $\delta$  is called the “Koszul-Tate differential” because it is associated with the Koszul-Tate resolution of the algebra of functions on the stationary surface (first step),

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# Koszul-Tate differential

The antifield-BRST approach to (gauge) field theories: an overview

Marc Henneaux

Introduction

BRST differential in Yang-Mills theory

Antifields and Koszul-Tate resolution

Homological perturbation theory

Local functionals and characteristic complex

Conclusions

The differential  $\delta$  is called the “Koszul-Tate differential” because it is associated with the Koszul-Tate resolution of the algebra of functions on the stationary surface (first step),

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We shall for this reason only focus here on the Koszul-Tate differential  $\delta$ .



# Koszul-Tate differential

## The antifield-BRST approach to (gauge) field theories: an overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

**Antifields and  
Koszul-Tate  
resolution**

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

# Koszul-Tate differential

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

**Antifields and  
Koszul-Tate  
resolution**

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

The algebra  $C^\infty(\Sigma)$  of smooth functions on the stationary surface can be viewed as the quotient of the algebra  $C^\infty(\mathcal{J})$  of smooth functions of the histories by the ideal  $\mathcal{N}$  of functions that vanish on  $\Sigma$ .

# Koszul-Tate differential

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

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# Koszul-Tate differential

The antifield-BRST approach to (gauge) field theories: an overview

Marc Henneaux

Introduction

BRST differential in Yang-Mills theory

Antifields and Koszul-Tate resolution

Homological perturbation theory

Local functionals and characteristic complex

Conclusions

The algebra  $C^\infty(\Sigma)$  of smooth functions on the stationary surface can be viewed as the quotient of the algebra  $C^\infty(J)$  of smooth functions of the histories by the ideal  $\mathcal{N}$  of functions that vanish on  $\Sigma$ .

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# Koszul-Tate differential

The antifield-BRST approach to (gauge) field theories: an overview

Marc Henneaux

Introduction

BRST differential in Yang-Mills theory

Antifields and Koszul-Tate resolution

Homological perturbation theory

Local functionals and characteristic complex

Conclusions

The algebra  $C^\infty(\Sigma)$  of smooth functions on the stationary surface can be viewed as the quotient of the algebra  $C^\infty(J)$  of smooth functions of the histories by the ideal  $\mathcal{N}$  of functions that vanish on  $\Sigma$ .

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# Koszul-Tate differential

The antifield-BRST approach to (gauge) field theories: an overview

Marc Henneaux

Introduction

BRST differential in Yang-Mills theory

Antifields and Koszul-Tate resolution

Homological perturbation theory

Local functionals and characteristic complex

Conclusions

The algebra  $C^\infty(\Sigma)$  of smooth functions on the stationary surface can be viewed as the quotient of the algebra  $C^\infty(J)$  of smooth functions of the histories by the ideal  $\mathcal{N}$  of functions that vanish on  $\Sigma$ .

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But this is exactly equivalent to  $f = \delta h$

# Koszul-Tate differential

The antifield-BRST approach to (gauge) field theories: an overview

Marc Henneaux

Introduction

BRST differential in Yang-Mills theory

Antifields and Koszul-Tate resolution

Homological perturbation theory

Local functionals and characteristic complex

Conclusions

The algebra  $C^\infty(\Sigma)$  of smooth functions on the stationary surface can be viewed as the quotient of the algebra  $C^\infty(J)$  of smooth functions of the histories by the ideal  $\mathcal{N}$  of functions that vanish on  $\Sigma$ .

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for some smooth coefficients  $k$ 's.

But this is exactly equivalent to  $f = \delta h$

with

$$h = k_\mu^a A_a^{*\mu} + k_\mu^{a\rho} \partial_\rho A_a^{*\mu} + k_\mu^{a\rho\sigma} \partial_\sigma \partial_\rho A_a^{*\mu} + \dots$$

# Koszul-Tate differential

## The antifield-BRST approach to (gauge) field theories: an overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

**Antifields and  
Koszul-Tate  
resolution**

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions



# Koszul-Tate differential

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

**Antifields and  
Koszul-Tate  
resolution**

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

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# Koszul-Tate differential

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

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# Koszul-Tate differential

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

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and therefore  $H_0(\delta) = C^\infty(\Sigma)$ .

Is there (co)homology at other values of the antifield number?

# Koszul-Tate differential

The antifield-BRST approach to (gauge) field theories: an overview

Marc Henneaux

Introduction

BRST differential in Yang-Mills theory

Antifields and Koszul-Tate resolution

Homological perturbation theory

Local functionals and characteristic complex

Conclusions

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Is there (co)homology at other values of the antifield number?

At antifield number 1, one finds that  $D_\mu A_a^{*\mu}$  is a cycle,  $\delta D_\mu A_a^{*\mu} = 0$  because of the Noether identity  $D_\mu D_\nu F_a^{\mu\nu} = 0$ .

# Koszul-Tate differential

The antifield-BRST approach to (gauge) field theories: an overview

Marc Henneaux

Introduction

BRST differential in Yang-Mills theory

Antifields and Koszul-Tate resolution

Homological perturbation theory

Local functionals and characteristic complex

Conclusions

Thus  $(\text{Im}\delta)_0 = \mathcal{N}$

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# Koszul-Tate differential

The antifield-BRST approach to (gauge) field theories: an overview

Marc Henneaux

Introduction

BRST differential in Yang-Mills theory

Antifields and Koszul-Tate resolution

Homological perturbation theory

Local functionals and characteristic complex

Conclusions

Thus  $(\text{Im}\delta)_0 = \mathcal{N}$

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# Koszul-Tate differential

The antifield-BRST approach to (gauge) field theories: an overview

Marc Henneaux

Introduction

BRST differential in Yang-Mills theory

Antifields and Koszul-Tate resolution

Homological perturbation theory

Local functionals and characteristic complex

Conclusions

Thus  $(\text{Im}\delta)_0 = \mathcal{N}$

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Indeed,

$$D_\mu A_a^{*\mu} = \delta C_a^*.$$

# Koszul-Tate differential

## The antifield-BRST approach to (gauge) field theories: an overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

**Antifields and  
Koszul-Tate  
resolution**

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions



# Koszul-Tate differential

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

**Antifields and  
Koszul-Tate  
resolution**

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

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# Koszul-Tate differential

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

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# Koszul-Tate differential

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

One can show that similarly,

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(If the gauge transformations were reducible, one would need “ghosts of ghosts” and on the Koszul-Tate side, “antifields for antifields”.)

# Koszul-Tate differential

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

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# Koszul-Tate differential

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

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Thus, the Koszul-Tate complex provides a resolution of the algebra  $C^\infty(\Sigma)$  of smooth functions on the stationary surface.

(If one includes the ghosts, one gets  $C^\infty(\Sigma) \otimes \Lambda(C^a, \partial_\mu C^a, \dots)$ .)

# Beyond Yang-Mills

## **The antifield-BRST approach to (gauge) field theories: an overview**

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

**Homological  
perturbation  
theory**

Local functionals  
and characteristic  
complex

Conclusions

# Beyond Yang-Mills

**The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview**

Marc Henneaux

The recognition of the antifields as related to a resolution of the stationary surface is key to the formulation of BRST theory beyond Yang-Mills.

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

**Homological  
perturbation  
theory**

Local functionals  
and characteristic  
complex

Conclusions

# Beyond Yang-Mills

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

The recognition of the antifields as related to a resolution of the stationary surface is key to the formulation of BRST theory beyond Yang-Mills.

(1) When the gauge transformations are reducible, one needs ghosts of ghosts and their conjugate antifields to maintain the resolution property.



# Beyond Yang-Mills

The antifield-BRST approach to (gauge) field theories: an overview

Marc Henneaux

Introduction

BRST differential in Yang-Mills theory

Antifields and Koszul-Tate resolution

Homological perturbation theory

Local functionals and characteristic complex

Conclusions

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(1) When the gauge transformations are reducible, one needs ghosts of ghosts and their conjugate antifields to maintain the resolution property.

(2) When the gauge transformations are “open” (on-shell closure only), the construction is more elaborate because  $\gamma^2 \neq 0$ , but  $\gamma^2 \approx 0$  (only on-shell). This requires additional terms in  $s$ ,

$$s = \delta + \gamma + s_1 + s_2 + \dots$$

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# Beyond Yang-Mills

The antifield-BRST approach to (gauge) field theories: an overview

Marc Henneaux

Introduction

BRST differential in Yang-Mills theory

Antifields and Koszul-Tate resolution

Homological perturbation theory

Local functionals and characteristic complex

Conclusions

The recognition of the antifields as related to a resolution of the stationary surface is key to the formulation of BRST theory beyond Yang-Mills.

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This is the Batalin-Vilkovisky construction, which works because the Koszul-Tate complex is a resolution.

# Homological perturbation theory

## The antifield-BRST approach to (gauge) field theories: an overview

Marc Henneaux

Introduction

BRST differential in Yang-Mills theory

Antifields and Koszul-Tate resolution

**Homological perturbation theory**

Local functionals and characteristic complex

Conclusions

# Homological perturbation theory

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

**Homological  
perturbation  
theory**

Local functionals  
and characteristic  
complex

Conclusions

To give an idea :

# Homological perturbation theory

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

**Homological  
perturbation  
theory**

Local functionals  
and characteristic  
complex

Conclusions

To give an idea :

$$(\delta + \gamma + \dots)^2 = \delta^2 + (\delta\gamma + \gamma\delta) + (\gamma^2) + \dots,$$

# Homological perturbation theory

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

**Homological  
perturbation  
theory**

Local functionals  
and characteristic  
complex

Conclusions

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# Homological perturbation theory

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

**Homological  
perturbation  
theory**

Local functionals  
and characteristic  
complex

Conclusions

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# Homological perturbation theory

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

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and therefore,

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# Homological perturbation theory

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

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# Homological perturbation theory

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

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(Homological perturbation theory).

# $p$ -form gauge fields, ghosts of ghosts

## The antifield-BRST approach to (gauge) field theories: an overview

Marc Henneaux

Introduction

BRST differential in Yang-Mills theory

Antifields and Koszul-Tate resolution

**Homological perturbation theory**

Local functionals and characteristic complex

Conclusions

# $p$ -form gauge fields, ghosts of ghosts

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

**Homological  
perturbation  
theory**

Local functionals  
and characteristic  
complex

Conclusions

Acyclicity of  $\delta$  is essential.

# $p$ -form gauge fields, ghosts of ghosts

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

**Homological  
perturbation  
theory**

Local functionals  
and characteristic  
complex

Conclusions

Acyclicity of  $\delta$  is essential.

If gauge transformations are reducible,  $\delta$  as defined above is not acyclic and one needs more antifields to recover this property :

# $p$ -form gauge fields, ghosts of ghosts

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

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$$B_{\mu\nu}, C_\mu, \gamma B_{\mu\nu} = \partial_{[\mu} C_{\nu]}, \delta C^{*\mu} = \partial_\nu B^{*\mu\nu}, \delta \partial_\mu C^{*\mu} = 0$$

# $p$ -form gauge fields, ghosts of ghosts

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

Acyclicity of  $\delta$  is essential.

If gauge transformations are reducible,  $\delta$  as defined above is not acyclic and one needs more antifields to recover this property :

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# $p$ -form gauge fields, ghosts of ghosts

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

Acyclicity of  $\delta$  is essential.

If gauge transformations are reducible,  $\delta$  as defined above is not acyclic and one needs more antifields to recover this property :

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# $p$ -form gauge fields, ghosts of ghosts

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

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On the ghost side, needs conjugate “ghost of ghosts”  $C$  with ghost number  $+2$  and such that  $\gamma C_\mu = \partial_\mu C$ .

# $p$ -form gauge fields, ghosts of ghosts

The antifield-BRST approach to (gauge) field theories: an overview

Marc Henneaux

Introduction

BRST differential in Yang-Mills theory

Antifields and Koszul-Tate resolution

Homological perturbation theory

Local functionals and characteristic complex

Conclusions

Acyclicity of  $\delta$  is essential.

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Procedure works and corresponding term in the solution of the master equation is  $\sim \int d^n x C^{*\mu} \partial_\mu C$ .

# Master equation

## The antifield-BRST approach to (gauge) field theories: an overview

Marc Henneaux

Introduction

BRST differential in Yang-Mills theory

Antifields and Koszul-Tate resolution

**Homological perturbation theory**

Local functionals and characteristic complex

Conclusions

# Master equation

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

**Homological  
perturbation  
theory**

Local functionals  
and characteristic  
complex

Conclusions

In terms of the solution  $S$  of the master equation  $(S, S) = 0$ ,

$$S = S_0 + S_1 + S_2 + \dots$$

# Master equation

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

**Homological  
perturbation  
theory**

Local functionals  
and characteristic  
complex

Conclusions

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$$(S_0, S_0) = 0, \quad (S_0, S_1) = 0, \quad 2(S_0, S_2) + (S_1, S_1) = 0, \dots$$

# Master equation

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

In terms of the solution  $S$  of the master equation  $(S, S) = 0$ ,

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# Master equation

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

In terms of the solution  $S$  of the master equation  $(S, S) = 0$ ,

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For instance,  $(S_0, S_1) = 0$  implies  $(S_0, (S_1, S_1)) = 0$  from which one infers the existence of  $S_2$  using acyclicity.

# Master equation

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

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Etc



# Extension

## **The antifield-BRST approach to (gauge) field theories: an overview**

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

**Homological  
perturbation  
theory**

Local functionals  
and characteristic  
complex

Conclusions

# Extension

## The antifield-BRST approach to (gauge) field theories: an overview

Marc Henneaux

Introduction

BRST differential in Yang-Mills theory

Antifields and Koszul-Tate resolution

**Homological perturbation theory**

Local functionals and characteristic complex

Conclusions

The construction can be carried through even if the equations do not derive from a variational principle.

# Extension

## The antifield-BRST approach to (gauge) field theories: an overview

Marc Henneaux

Introduction

BRST differential in Yang-Mills theory

Antifields and Koszul-Tate resolution

**Homological perturbation theory**

Local functionals and characteristic complex

Conclusions

The construction can be carried through even if the equations do not derive from a variational principle.

In that case, however, there is in general no matching between fields and antifields

# Extension

## The antifield-BRST approach to (gauge) field theories: an overview

Marc Henneaux

Introduction

BRST differential in Yang-Mills theory

Antifields and Koszul-Tate resolution

Homological perturbation theory

Local functionals and characteristic complex

Conclusions

The construction can be carried through even if the equations do not derive from a variational principle.

In that case, however, there is in general no matching between fields and antifields

and no natural antibracket.

# Extension

## The antifield-BRST approach to (gauge) field theories: an overview

Marc Henneaux

Introduction

BRST differential in Yang-Mills theory

Antifields and Koszul-Tate resolution

Homological perturbation theory

Local functionals and characteristic complex

Conclusions

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# Extension

## The antifield-BRST approach to (gauge) field theories: an overview

Marc Henneaux

Introduction

BRST differential in Yang-Mills theory

Antifields and Koszul-Tate resolution

Homological perturbation theory

Local functionals and characteristic complex

Conclusions

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In that case, however, there is in general no matching between fields and antifields

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The BRST differential  $s = \delta + \gamma + \dots$  can be constructed as before but is not generated in the antibracket through the solution  $S$  of the master equation.

# Local Functionals

## The antifield-BRST approach to (gauge) field theories: an overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

**Local functionals  
and characteristic  
complex**

Conclusions

# Local Functionals

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

**Local functionals  
and characteristic  
complex**

Conclusions

Also quite relevant to physics is the BRST cohomology in the space of local functionals.



# Local Functionals

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

**Local functionals  
and characteristic  
complex**

Conclusions

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**A local functional is the integral of a local  $n$ -form**

# Local Functionals

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

**Local functionals  
and characteristic  
complex**

Conclusions

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**A local functional is the integral of a local  $n$ -form**

A local  $p$ -form is a  $p$ -form with coefficients that are local functions.

# Local Functionals

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

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A local functional is the integral of a local  $n$ -form

A local  $p$ -form is a  $p$ -form with coefficients that are local functions.

So a local functional is

$$F = \int \omega, \quad \omega = f d^n x,$$

where  $f$  is a local function.

# BRST cohomology in the space of local functionals

## The antifield-BRST approach to (gauge) field theories: an overview

Marc Henneaux

Introduction

BRST differential in Yang-Mills theory

Antifields and Koszul-Tate resolution

Homological perturbation theory

**Local functionals and characteristic complex**

Conclusions

# BRST cohomology in the space of local functionals

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

**Local functionals  
and characteristic  
complex**

Conclusions

The cocycles and coboundary conditions for local functionals  
read

# BRST cohomology in the space of local functionals

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

**Local functionals  
and characteristic  
complex**

Conclusions

The cocycles and coboundary conditions for local functionals read

$$s\omega + da = 0 \quad \text{and} \quad \omega = s\psi + db$$

# BRST cohomology in the space of local functionals

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

**Local functionals  
and characteristic  
complex**

Conclusions

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$$s\omega + da = 0 \quad \text{and} \quad \omega = s\psi + db$$

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# BRST cohomology in the space of local functionals

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

**Local functionals  
and characteristic  
complex**

Conclusions

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# BRST cohomology in the space of local functionals

The antifield-BRST approach to (gauge) field theories: an overview

Marc Henneaux

Introduction

BRST differential in Yang-Mills theory

Antifields and Koszul-Tate resolution

Homological perturbation theory

Local functionals and characteristic complex

Conclusions

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in terms of the integrands since

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This defines the mod- $d$  cohomology  $H^m(s|d)$ .

# $H_m(\delta|d)$

## The antifield-BRST approach to (gauge) field theories: an overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

**Local functionals  
and characteristic  
complex**

Conclusions

$$H_m(\delta|d)$$

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

**Local functionals  
and characteristic  
complex**

Conclusions

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it turns out not to be true that  $H_m(s|d) = 0$  for  $m > 0$ .

For instance an abelian ghost  $C^*$  fulfills  $\delta C^* = \partial_\mu A^{*\mu}$  and defines a non-trivial element of  $H_2(\delta|d)$ .

# Characteristic cohomology

## The antifield-BRST approach to (gauge) field theories: an overview

Marc Henneaux

Introduction

BRST differential in Yang-Mills theory

Antifields and Koszul-Tate resolution

Homological perturbation theory

**Local functionals and characteristic complex**

Conclusions



# Characteristic cohomology

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

**Local functionals  
and characteristic  
complex**

Conclusions

The characteristic cohomology in form-degree  $k$  is defined

# Characteristic cohomology

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

**Local functionals  
and characteristic  
complex**

Conclusions

The characteristic cohomology in form-degree  $k$  is defined to be the space of  $k$ -forms that are closed on-shell

# Characteristic cohomology

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

**Local functionals  
and characteristic  
complex**

Conclusions

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# Characteristic cohomology

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

**Local functionals  
and characteristic  
complex**

Conclusions

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# Characteristic cohomology

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

The characteristic cohomology in form-degree  $k$  is defined to be the space of  $k$ -forms that are closed on-shell modulo the  $k$ -forms that are exact on-shell,

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# Characteristic cohomology

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

The characteristic cohomology in form-degree  $k$  is defined to be the space of  $k$ -forms that are closed on-shell modulo the  $k$ -form that are exact on-shell,

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# Characteristic cohomology

The antifield-BRST approach to (gauge) field theories: an overview

Marc Henneaux

Introduction

BRST differential in Yang-Mills theory

Antifields and Koszul-Tate resolution

Homological perturbation theory

Local functionals and characteristic complex

Conclusions

The characteristic cohomology in form-degree  $k$  is defined to be the space of  $k$ -forms that are closed on-shell modulo the  $k$ -form that are exact on-shell,

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Using the Koszul-Tate differential, one easily sees that the characteristic cohomology is just  $H_0^k(d|\delta)$ .

Reading the cocycle condition  $\delta\omega + da = 0$  in both directions, one easily proves the isomorphisms

$$H_j^i(\delta|d) \simeq H_{j-1}^{i-1}(d|\delta), \quad i, j > 1, \quad (i, j) \neq (1, 1); \quad H_1^1(\delta|d) \simeq \frac{H_0^0(d|\delta)}{\mathbb{R}}$$

# Results

## **The antifield-BRST approach to (gauge) field theories: an overview**

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

**Local functionals  
and characteristic  
complex**

Conclusions



# Results

## The antifield-BRST approach to (gauge) field theories: an overview

Marc Henneaux

Introduction

BRST differential in Yang-Mills theory

Antifields and Koszul-Tate resolution

Homological perturbation theory

**Local functionals and characteristic complex**

Conclusions

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# Results

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

In particular,  $H_1^n(\delta|\mathcal{d}) \simeq H_0^{n-1}(\mathcal{d}|\delta)$  is just a cohomological reformulation of the Noether theorem.

$H_2^n(\delta|\mathcal{d}) \simeq H_0^{n-2}(\mathcal{d}|\delta)$  relates  $*F$  (which is a  $(n-2)$ -form closed on-shell in the abelian case) to  $c^* d^n x$ .

# Results

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

In particular,  $H_1^n(\delta|d) \simeq H_0^{n-1}(d|\delta)$  is just a cohomological reformulation of the Noether theorem.

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One can compute explicitly the cohomologies  $H(\delta|d)$  and  $H(d|\delta)$

# Results

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

In particular,  $H_1^n(\delta|d) \simeq H_0^{n-1}(d|\delta)$  is just a cohomological reformulation of the Noether theorem.

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# Results

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

In particular,  $H_1^n(\delta|d) \simeq H_0^{n-1}(d|\delta)$  is just a cohomological reformulation of the Noether theorem.

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# Results

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

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Once this is done, one can compute explicitly the BRST cohomology  $H(s|d)$  (BRST cohomology in the space of local functionals).

The understanding that  $s$  involves  $\delta$  – and the corresponding spectral sequence – is crucial for this purpose.

# Conclusions and comments

## The antifield-BRST approach to (gauge) field theories: an overview

Marc Henneaux

Introduction

BRST differential in Yang-Mills theory

Antifields and Koszul-Tate resolution

Homological perturbation theory

Local functionals and characteristic complex

**Conclusions**

# Conclusions and comments

## The antifield-BRST approach to (gauge) field theories: an overview

Marc Henneaux

Introduction

BRST differential in Yang-Mills theory

Antifields and Koszul-Tate resolution

Homological perturbation theory

Local functionals and characteristic complex

**Conclusions**

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# Conclusions and comments

## The antifield-BRST approach to (gauge) field theories: an overview

Marc Henneaux

Introduction

BRST differential in Yang-Mills theory

Antifields and Koszul-Tate resolution

Homological perturbation theory

Local functionals and characteristic complex

Conclusions

Antifields were originally introduced by Zinn-Justin as sources coupled to the BRST variations of the fields.

**A different interpretation of the antifields can be developed.**

# Conclusions and comments

## The antifield-BRST approach to (gauge) field theories: an overview

Marc Henneaux

Introduction

BRST differential in Yang-Mills theory

Antifields and Koszul-Tate resolution

Homological perturbation theory

Local functionals and characteristic complex

Conclusions

Antifields were originally introduced by Zinn-Justin as sources coupled to the BRST variations of the fields.

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The antifields can indeed also be viewed as the generators of the Koszul-Tate “resolution” that implements the equations of motion in cohomology.

# Conclusions and comments

## The antifield-BRST approach to (gauge) field theories: an overview

Marc Henneaux

Introduction

BRST differential in Yang-Mills theory

Antifields and Koszul-Tate resolution

Homological perturbation theory

Local functionals and characteristic complex

Conclusions

Antifields were originally introduced by Zinn-Justin as sources coupled to the BRST variations of the fields.

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**This point of view turns out to be crucial for computing the BRST cohomology.**

# Conclusions and comments

## The antifield-BRST approach to (gauge) field theories: an overview

Marc Henneaux

Introduction

BRST differential in Yang-Mills theory

Antifields and Koszul-Tate resolution

Homological perturbation theory

Local functionals and characteristic complex

**Conclusions**

# Conclusions and comments

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

**Conclusions**

# THANK YOU!

**The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview**

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

**Conclusions**

**The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview**

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

**Conclusions**

# Backup Material

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

## REFERENCES



# Homological perturbation theory (“strong homotopy Lie algebras $L_\infty$ ”)

## The antifield-BRST approach to (gauge) field theories: an overview

Marc Henneaux

Introduction

BRST differential in Yang-Mills theory

Antifields and Koszul-Tate resolution

Homological perturbation theory

Local functionals and characteristic complex

Conclusions

# Homological perturbation theory (“strong homotopy Lie algebras $L_\infty$ ”)

The antifield-BRST approach to (gauge) field theories: an overview

Marc Henneaux

Introduction

BRST differential in Yang-Mills theory

Antifields and Koszul-Tate resolution

Homological perturbation theory

Local functionals and characteristic complex

Conclusions

I.A. Batalin and G.A. Vilkovisky, *Gauge Algebra and Quantization*, Phys.Lett. **102B** (1981) 27-31 ;  
*Quantization of Gauge Theories with Linearly Dependent Generators*, Phys.Rev. **D28** (1983) 2567-2582, Erratum : Phys.Rev. **D30** (1984) 508

# Homological perturbation theory (“strong homotopy Lie algebras $L_\infty$ ”)

The antifield-BRST approach to (gauge) field theories: an overview

Marc Henneaux

Introduction

BRST differential in Yang-Mills theory

Antifields and Koszul-Tate resolution

Homological perturbation theory

Local functionals and characteristic complex

Conclusions

I.A. Batalin and G.A. Vilkovisky, *Gauge Algebra and Quantization*, Phys.Lett. **102B** (1981) 27-31 ;

*Quantization of Gauge Theories with Linearly Dependent Generators*, Phys.Rev. **D28** (1983) 2567-2582, Erratum : Phys.Rev. **D30** (1984) 508

J. Fisch and M. Henneaux, *Homological Perturbation Theory and the Algebraic Structure of the Antifield - Antibracket Formalism for Gauge Theories*, Commun.Math.Phys. **128** (1990) 627 ;

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# Homological perturbation theory (“strong homotopy Lie algebras $L_\infty$ ”)

The antifield-BRST approach to (gauge) field theories: an overview

Marc Henneaux

Introduction

BRST differential in Yang-Mills theory

Antifields and Koszul-Tate resolution

Homological perturbation theory

Local functionals and characteristic complex

Conclusions

I.A. Batalin and G.A. Vilkovisky, *Gauge Algebra and Quantization*, Phys.Lett. **102B** (1981) 27-31 ;

*Quantization of Gauge Theories with Linearly Dependent Generators*, Phys.Rev. **D28** (1983) 2567-2582, Erratum : Phys.Rev. **D30** (1984) 508

J. Fisch and M. Henneaux, *Homological Perturbation Theory and the Algebraic Structure of the Antifield - Antibracket Formalism for Gauge Theories*, Commun.Math.Phys. **128** (1990) 627 ;

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(Hamiltonian precursor in :

# Homological perturbation theory (“strong homotopy Lie algebras $L_\infty$ ”)

The antifield-BRST approach to (gauge) field theories: an overview

Marc Henneaux

Introduction

BRST differential in Yang-Mills theory

Antifields and Koszul-Tate resolution

Homological perturbation theory

Local functionals and characteristic complex

Conclusions

I.A. Batalin and G.A. Vilkovisky, *Gauge Algebra and Quantization*, Phys.Lett. **102B** (1981) 27-31 ;

*Quantization of Gauge Theories with Linearly Dependent Generators*, Phys.Rev. **D28** (1983) 2567-2582, Erratum : Phys.Rev. **D30** (1984) 508

J. Fisch and M. Henneaux, *Homological Perturbation Theory and the Algebraic Structure of the Antifield - Antibracket Formalism for Gauge Theories*, Commun.Math.Phys. **128** (1990) 627 ;

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# Homological perturbation theory (“strong homotopy Lie algebras $L_\infty$ ”)

The antifield-BRST approach to (gauge) field theories: an overview

Marc Henneaux

Introduction

BRST differential in Yang-Mills theory

Antifields and Koszul-Tate resolution

Homological perturbation theory

Local functionals and characteristic complex

Conclusions

I.A. Batalin and G.A. Vilkovisky, *Gauge Algebra and Quantization*, Phys.Lett. **102B** (1981) 27-31 ;

*Quantization of Gauge Theories with Linearly Dependent Generators*, Phys.Rev. **D28** (1983) 2567-2582, Erratum : Phys.Rev. **D30** (1984) 508

J. Fisch and M. Henneaux, *Homological Perturbation Theory and the Algebraic Structure of the Antifield - Antibracket Formalism for Gauge Theories*, Commun.Math.Phys. **128** (1990) 627 ;

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# Cohomology

## The antifield-BRST approach to (gauge) field theories: an overview

Marc Henneaux

Introduction

BRST differential in Yang-Mills theory

Antifields and Koszul-Tate resolution

Homological perturbation theory

Local functionals and characteristic complex

Conclusions

# Cohomology

## The antifield-BRST approach to (gauge) field theories: an overview

Marc Henneaux

Introduction

BRST differential in Yang-Mills theory

Antifields and Koszul-Tate resolution

Homological perturbation theory

Local functionals and characteristic complex

Conclusions

The connection between the BRST cohomology and the characteristic cohomology was made in



# Cohomology

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

The connection between the BRST cohomology and the characteristic cohomology was made in

G. Barnich, F. Brandt, MH, *Local BRST cohomology in the antifield formalism. 1. General theorems*, Commun.Math.Phys. **174** (1995) 57-92

# Cohomology

The  
antifield-BRST  
approach to  
(gauge) field  
theories: an  
overview

Marc Henneaux

Introduction

BRST differential  
in Yang-Mills  
theory

Antifields and  
Koszul-Tate  
resolution

Homological  
perturbation  
theory

Local functionals  
and characteristic  
complex

Conclusions

The connection between the BRST cohomology and the characteristic cohomology was made in

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